

Rec 1

Monday, October 2, 2023 01:07

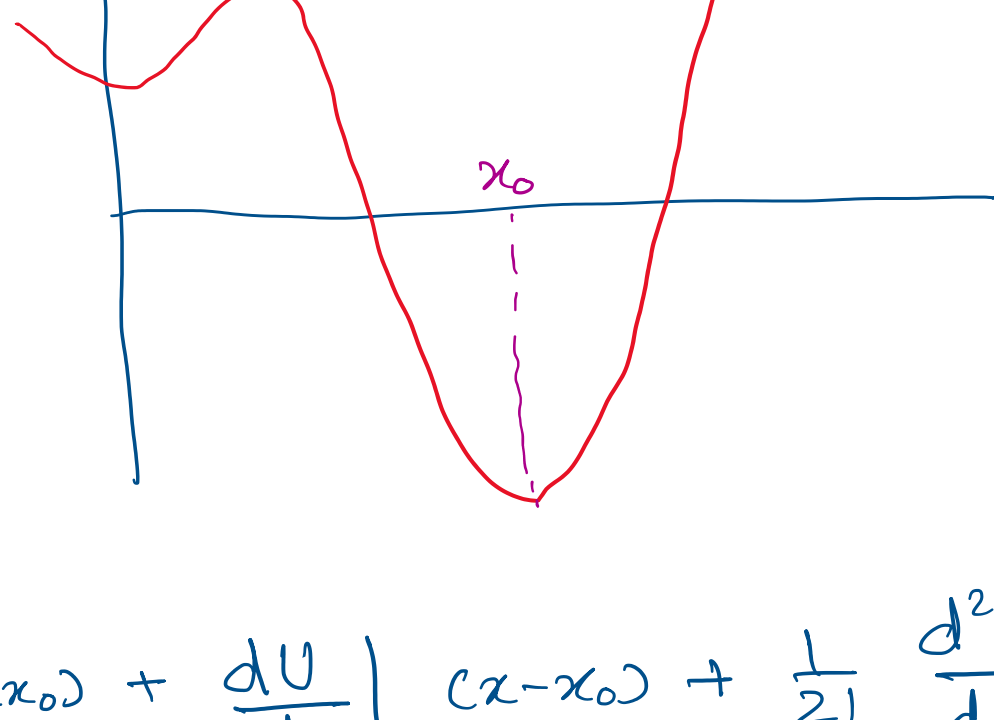
You have learnt a lot about SHO's this week. One of the reasons we focus on this topic is because it's one of the few things we know how to solve.

"The career of a young theoretical physicist is treating HO in increasing levels of abstraction" - Sydney Coleman

So, instead of learning new laws, we will focus on applying what we know to increasingly complex systems.

Let's start with understanding where SHO's occur.

For any system with conservative forces



$$U(x) = U(x_0) + \left. \frac{dU}{dx} \right|_{x_0} (x-x_0) + \frac{1}{2!} \left. \frac{d^2U}{dx^2} \right|_{x_0} (x-x_0)^2 + \dots$$

(Taylor expanded about equilibrium)

Now, using $F = -\frac{dU}{dx}$

$$F = 0 - \left. \frac{dU}{dx} \right|_{x=x_0} - \left. \frac{d^2U}{dx^2} \right|_{x=x_0} (x-x_0) - \dots$$

\downarrow
 ≈ 0

$$\therefore F = - \left. \frac{d^2U}{dx^2} \right|_{x_0} (x-x_0)$$

$$\& U = U(x_0) + \frac{1}{2!} \left. \frac{d^2U}{dx^2} \right|_{x_0} (x-x_0)^2$$

1) Let's apply this to the case of diatomic molecules

$$U(x) = -\frac{a}{x^6} + \frac{b}{x^{12}} ; a, b > 0$$

↖ attractive force
↘ repulsive force

i) is this SHO?

Not yet

ii) Find the minima

$$F(x) = 0 = -\frac{6a}{x_0^7} + \frac{12b}{x_0^{13}}$$

$$\Rightarrow x_0 = (2b/a)^{1/6}$$

iii) Find the eqns of motion

$$F(x) = -\frac{6a}{(x_0+\Delta x)^7} + \frac{12b}{(x_0+\Delta x)^{13}}$$

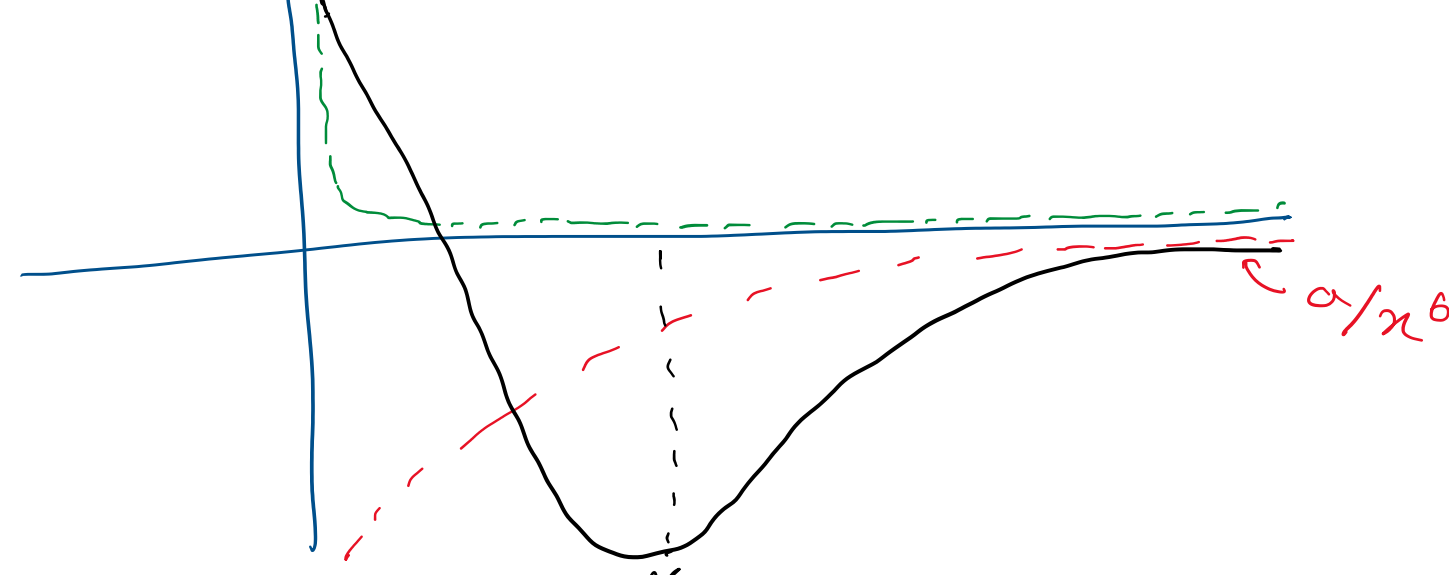
$$= -\frac{6a}{x_0^7} \left(1 + \frac{\Delta x}{x_0}\right)^{-7} + \frac{12b}{x_0^{13}} \left(1 + \frac{\Delta x}{x_0}\right)^{-13}$$

↖ ← are equal

$$= \alpha \left[-1 + \frac{7\Delta x}{x_0} + 1 - \frac{13\Delta x}{x_0} \right]$$

$$= -\frac{6\alpha}{x_0} \Delta x$$

$$= -\frac{36a}{x_0^8} \Delta x$$



2) Let's now solve the eqns of motion using conservation of energy

$$E = \frac{1}{2} mv^2 + mgl(1 - \cos\theta)$$

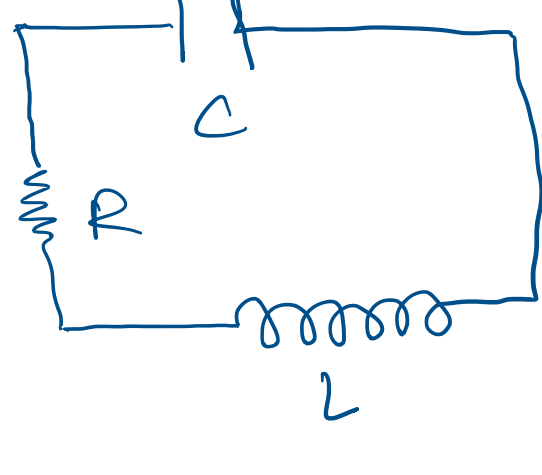
$$\Rightarrow E = \frac{1}{2} ml^2 \dot{\theta}^2 + mgl(1 - \cos\theta)$$

$$\Rightarrow \frac{dE}{dt} = 0 = ml^2 \dot{\theta} \ddot{\theta} + mgl \sin\theta \dot{\theta}$$

using small angle approximation,

$$\ddot{\theta} = -\frac{g}{l} \theta$$

3) Let's make life a bit more complicated.



$$\frac{q}{c} - L \frac{dI}{dt} - IR = 0$$

$$\Rightarrow L \frac{dI}{dt} - \frac{q}{c} + IR = 0$$

$$\boxed{I = -dq/dt} \Rightarrow L \frac{d^2q}{dt^2} + R \frac{dq}{dt} + \frac{q}{c} = 0$$

$$\Rightarrow \ddot{q} + \frac{R}{L} \dot{q} + \frac{q}{LC} = 0$$

Let's solve the generic eqn instead.

$$\ddot{z} + \gamma \dot{z} + \omega_0^2 z = 0$$

$$\text{Ansatz} \rightarrow z(t) = Ce^{\alpha t}$$

$$\alpha^2 + \gamma\alpha + \omega_0^2 = 0$$

$$\alpha = \frac{-\gamma \pm \sqrt{\gamma^2 - 4\omega_0^2}}{2}$$

$$\therefore z(t) = C_1 e^{\alpha_1 t} + C_2 e^{\alpha_2 t}$$

$$= e^{-\gamma/2 t} \left[C_1 e^{+\frac{\sqrt{\gamma^2 - 4\omega_0^2}}{2} t} + C_2 e^{-\frac{\sqrt{\gamma^2 - 4\omega_0^2}}{2} t} \right]$$

always dies out

4) How about a driven oscillator,

$$\ddot{x} + \gamma \dot{x} + \omega_0^2 x = \frac{F_0}{m} \cos(\omega t)$$

so how do you find the right ansatz

$$\ddot{y} - 4\dot{y} - 12y = 3e^{st}$$

$$\text{Ansatz} \rightarrow Ae^{st}$$

$$\ddot{y} - 4\dot{y} - 12y = 2t^2 - t^2$$

$$\text{Ansatz} \rightarrow At^2 + Bt + C + D$$

$$\therefore \text{for } \ddot{z} + \gamma \dot{z} + \omega_0^2 z = \frac{F_0}{m} e^{i\omega t}$$

$$\therefore z_p(t) = A e^{i(\omega t - \delta)}$$

\downarrow
ansatz

$$\therefore A e^{i(\omega t - \delta)} [-\omega^2 + i\gamma\omega + \omega_0^2] = \frac{F_0}{m} e^{i\omega t}$$

$$(\omega_0^2 - \omega^2)A + i\gamma\omega A = \frac{F_0}{m} e^{i\delta}$$

$$\frac{F_0}{m} \sin\delta = \gamma\omega A$$

$$\frac{F_0}{m} \cos\delta = (\omega_0^2 - \omega^2)A$$

$$\therefore A(\omega) = \frac{F_0/m}{\sqrt{(\omega_0^2 - \omega^2)^2 + \gamma^2\omega^2}} \quad \delta = \tan^{-1} \left(\frac{\gamma\omega}{\omega_0^2 - \omega^2} \right)$$